

Hamilton–Jacobi theory for contact systems: autonomous and non-autonomous

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Geometric Hamilton–Jacobi theory

- Consider a dynamical system characterized by $X \in \mathfrak{X}(M)$.
- Suppose that $\pi: M \rightarrow B$ is a vector bundle (e.g. $\pi_Q: T^*Q \rightarrow Q$).
- Idea: obtain a section $\gamma \in \Gamma(M)$ such that the following diagram commutes:

$$\begin{array}{ccc}
 M & \xrightarrow{X} & TM \\
 \left. \begin{array}{c} \nearrow \gamma \\ \downarrow \pi \\ \searrow \end{array} \right\} & & \left. \begin{array}{c} \downarrow T\pi \\ \nearrow T\gamma \end{array} \right\} \\
 B & \xrightarrow{X^\gamma} & TB
 \end{array}$$

- If σ is an integral curve of X^γ , then $\gamma \circ \sigma$ is an integral curve of X .

Cosymplectic and contact structures

Let M be a $(2n + 1)$ -dimensional manifold

Cosymplectic manifold (M, ω, τ) Contact manifold (M, η)

- ω closed 2-form
- τ closed 1-form
- $\tau \wedge \omega^n \neq 0$
- Reeb vector field \mathcal{R}_t :

$$\iota_{\mathcal{R}_t} \omega = 0, \quad \iota_{\mathcal{R}_t} \tau = 1$$

- Darboux coords. (t, q^i, p_i) :

$$\omega = dq^i \wedge dp_i, \quad \tau = dt, \quad \mathcal{R}_t = \frac{\partial}{\partial t}$$

- η 1-form

- $\eta \wedge d\eta^n \neq 0$

- Reeb vector field \mathcal{R}_t :

$$\iota_{\mathcal{R}_t} \eta = 1, \quad \iota_{\mathcal{R}_t} d\eta = 0$$

- Darboux coords. (q^i, p_i, z) :

$$\eta = dz - p_i dq^i, \quad \mathcal{R}_z = \frac{\partial}{\partial z}$$

Cocontact structures I

- Idea: a structure that combines the cosymplectic and contact ones.

Definition

A **cocontact manifold** is a triple (M, τ, η) where:

- 1 M is a $(2n + 2)$ -dimensional manifold,
- 2 τ and η are 1-forms,
- 3 $d\tau = 0$,
- 4 $\tau \wedge \eta \wedge (d\eta)^{\wedge n} \neq 0$.

Cocontact structures II

- Given a cocontact manifold (M, τ, η) , we have the **flat isomorphism**:

$$b: \mathfrak{X}(M) \rightarrow \Omega^1(M)$$

$$X \mapsto (\iota_X \tau)\tau + \iota_X d\eta + (\iota_X \eta)\eta$$

and its inverse $\sharp = b^{-1}$.

- Reeb vector fields:** $\mathcal{R}_t = b^{-1}(\tau)$, $\mathcal{R}_z = b^{-1}(\eta)$.
- Darboux coordinates (t, q^i, p_i, z) :

$$\tau = dt, \quad \eta = dz - p_i dq^i, \quad \mathcal{R}_t = \frac{\partial}{\partial t}, \quad \mathcal{R}_z = \frac{\partial}{\partial z}$$

Cocontact Hamiltonian systems

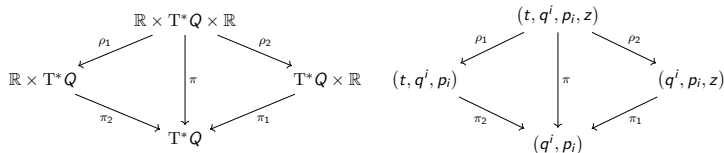
- Given a Hamiltonian function $H: M \rightarrow \mathbb{R}$, its **Hamiltonian vector field** is given by

$$\flat(X_H) = dH - (\mathcal{R}_z(H) + H)\eta + (1 - \mathcal{R}_t(H))\tau.$$

- In Darboux coordinates,

$$X_H = \frac{\partial}{\partial t} + \frac{\partial H}{\partial p_i} \frac{\partial}{\partial q^i} - \left(\frac{\partial H}{\partial q^i} + p_i \frac{\partial H}{\partial z} \right) \frac{\partial}{\partial p_i} + \left(p_i \frac{\partial H}{\partial p_i} - H \right) \frac{\partial}{\partial z}.$$

Canonical cocontact manifold



- Let Q be an n -dimensional manifold with local coordinates (q^i) .
- Let $\theta_0 = p_i dq^i$ be the canonical 1-form of T^*Q .
- Consider the 1-forms $\theta_Q = \pi^*\theta_0$ and $\eta_Q = dz - \theta_Q$ on $\mathbb{R} \times T^*Q \times \mathbb{R}$
- Then, (dt, η_Q) is a cocontact structure on $\mathbb{R} \times T^*Q \times \mathbb{R}$. The local expression of the 1-form η is

$$\eta_Q = dz - p_i dq^i.$$

The action-independent approach

- Let $(\mathbb{R} \times T^*Q \times \mathbb{R}, dt, \eta_Q, H)$ be a cocontact Hamiltonian system.
- Idea: obtain a section γ of $\pi_Q^t: \mathbb{R} \times T^*Q \times \mathbb{R} \rightarrow \mathbb{R} \times Q$ such that the following diagram commutes:

$$\begin{array}{ccc}
 \mathbb{R} \times T^*Q \times \mathbb{R} & \xrightarrow{X_H} & T(\mathbb{R} \times T^*Q \times \mathbb{R}) \\
 \gamma \left(\begin{array}{c} \uparrow \\ \downarrow \pi_Q^t \end{array} \right) & & \left(\begin{array}{c} \uparrow \\ \downarrow T\pi_Q^t \end{array} \right) T\gamma \\
 \mathbb{R} \times Q & \xrightarrow{X_H^\gamma} & T(\mathbb{R} \times Q)
 \end{array}$$

- Here $\pi_Q^t: (t, q^i, p_i, z) \mapsto (t, q^i)$.

Proposition

$\text{Im}\gamma(t, \cdot)$ is a Legendrian submanifold $\forall t \in \mathbb{R}$ (i.e., $\gamma^*\eta_Q = 0$) iff

$$\gamma(t, q) = j_t^1 S(t, q) := \left(t, q^i, \frac{\partial S}{\partial q^i}, S \right)$$

Theorem (Action-independent Hamilton–Jacobi Theorem)

Suppose that, $\forall t \in \mathbb{R}$, $\text{Im}\gamma(t, \cdot)$ is a Legendrian submanifold. Then, X_H^γ and X_H are γ -related iff

$$H \circ j_t^1 S + \frac{\partial S}{\partial t} = 0.$$

The function S is called a **generating function** for H .

Application: time-independent contact systems

- Let $(T^*Q \times \mathbb{R}, \eta_Q, H)$ be a contact Hamiltonian system.
- Consider the associated cocontact Hamiltonian system $(\mathbb{R} \times T^*Q \times \mathbb{R}, dt, \eta_Q, H \circ \rho_2)$, where $\rho_2: (t, q^i, p_i, z) \mapsto (q^i, p_i, z)$.
- Suppose that $S(t, q) = \alpha(q) + \beta(t)$.
- Then, the Hamilton–Jacobi equation is written as

$$H \circ j^1\alpha + \frac{\partial \beta}{\partial t} = 0,$$

- Notice that if S is time-independent (i.e., $\beta = 0$), the solutions of the HJ problem only cover the zero energy level!

Example (The free particle with linear dissipation)

- Consider the contact Hamiltonian system $(T^*\mathbb{R} \times \mathbb{R}, \eta_{\mathbb{R}}, H)$ with

$$H(q, p, z) = \frac{p^2}{2} + \delta z.$$

- Then, $S(t, q) = \lambda e^{-\delta t} - \frac{\delta}{2} q^2$ is a generating function for H .
- Now,

$$X_H^\gamma = \frac{\partial}{\partial t} + \frac{\partial H}{\partial p} \frac{\partial}{\partial q} \Big|_{\text{Im } \gamma} = \frac{\partial}{\partial t} - \delta q \frac{\partial}{\partial q},$$

whose integral curves are of the form $\sigma(t) = (t, q_0 e^{-\delta t})$.

- Thus, the integral curves of $X_H|_{\text{Im}(\gamma)}$ are given by

$$\gamma \circ \sigma(t) = (t, q_0 e^{-\delta t}, -\delta q_0 e^{-\delta t}, -\frac{\delta}{2} q_0^2 e^{-2\delta t} + \lambda e^{-\delta t}).$$

The action-dependent approach

- The previous approach has a drawback: complete solutions cannot be defined.
- Idea: consider sections of $\pi_Q^{t,z}: \mathbb{R} \times T^*Q \times \mathbb{R} \rightarrow \mathbb{R} \times Q \times \mathbb{R}$:

$$\begin{array}{ccc}
 \mathbb{R} \times T^*Q \times \mathbb{R} & \xrightarrow{X_H} & T(\mathbb{R} \times T^*Q \times \mathbb{R}) \\
 \gamma \left(\begin{array}{c} \uparrow \\ \downarrow \end{array} \right) \pi_Q^{t,z} & & T\pi_Q^{t,z} \left(\begin{array}{c} \uparrow \\ \downarrow \end{array} \right) T\gamma \\
 \mathbb{R} \times Q \times \mathbb{R} & \xrightarrow{X_H^\gamma} & T(\mathbb{R} \times Q \times \mathbb{R})
 \end{array}$$

Proposition

If (M, τ, η) is a cocontact manifold, then $(M, \Lambda, -\mathcal{R}_z)$ is a Jacobi manifold, where $\Lambda(\alpha, \beta) = -d\eta(\#\alpha, \#\beta)$.

- Recall that the orthogonal complement \mathcal{D}^\perp of a distribution $\mathcal{D} \subseteq TM$ is given by $\mathcal{D}^\perp = \Lambda(\mathcal{D}^\circ, \cdot)$.
- A submanifold N is said to be **coisotropic** if $TN^\perp \subseteq TN$.

- Let $d_Q f := \frac{\partial f}{\partial q^i} dq^i$ for $f \in C^\infty(\mathbb{R} \times Q \times \mathbb{R})$.

Theorem (Action-dependent Hamilton–Jacobi Theorem)

Let γ be a section of $\pi_Q^{t,s} : \mathbb{R} \times T^*Q \times \mathbb{R} \rightarrow \mathbb{R} \times Q \times \mathbb{R}$ such that $\text{Im } \gamma$ is a coisotropic submanifold. Then, X_H^γ and X_H are γ -related iff

$$d_Q(H \circ \gamma) + \mathcal{L}_{\mathcal{R}_z}(H \circ \gamma) \gamma + \mathcal{L}_{\mathcal{R}_t} \gamma = (H \circ \gamma) \mathcal{L}_{\mathcal{R}_z} \gamma.$$

Complete solutions I

Definition

A **complete solution of the Hamilton–Jacobi problem** is a local diffeomorphism $\Phi: \mathbb{R} \times Q \times \mathbb{R}^n \times \mathbb{R} \rightarrow \mathbb{R} \times T^*Q \times \mathbb{R}$ such that, for each $\lambda \in \mathbb{R}^n$,

$$\begin{aligned} \Phi_\lambda: \mathbb{R} \times Q \times \mathbb{R} &\longrightarrow \mathbb{R} \times T^*Q \times \mathbb{R} \\ (t, q^i, z) &\longmapsto \Phi(t, q^i, \lambda, z) \end{aligned}$$

is a solution of the action-dependent HJ problem.

Complete solutions II

- Let us define the functions $f_i = \pi_i \circ \alpha \circ \Phi^{-1}$ on $\mathbb{R} \times T^*Q \times \mathbb{R}$, so that the following diagram commutes:

$$\begin{array}{ccc}
 \mathbb{R} \times Q \times \mathbb{R} \times \mathbb{R}^n & \begin{array}{c} \xrightarrow{\Phi} \\ \xleftarrow{\Phi^{-1}} \end{array} & \mathbb{R} \times T^*Q \times \mathbb{R} \\
 \downarrow \alpha & & \downarrow f_i \\
 \mathbb{R}^n & \xrightarrow{\pi_i} & \mathbb{R}
 \end{array}$$

Theorem

For each $i \in \{1, \dots, n\}$, the function $f_i = \pi_i \circ \alpha \circ \Phi^{-1}$ is a constant of the motion. However, these functions are not necessarily in involution, i.e., $\{f_i, f_j\} \neq 0$.

Freely falling particle with linear dissipation I

- The Hamiltonian function $H : \mathbb{R} \times \mathbb{T}^*\mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}$ is given by

$$H(t, q, p, z) = \frac{p^2}{2m(t)} + m(t)gq + \frac{\gamma}{m(t)}z.$$

- We look for a conserved quantity f , i.e., $X_H(f) = 0$.
- For simplicity's sake, suppose that $f = f(t, p)$.
- Conserved quantity:

$$f(t, q, p, z) = e^{\int_1^t \frac{\gamma}{m(s)} ds} \left(p + ge^{-\int_1^t \frac{\gamma}{m(s)} ds} \int_1^t e^{-\int_1^u -\frac{\gamma}{m(s)} ds} m(u) du \right).$$

Freely falling particle with linear dissipation II

- We can thus express the momentum p as a function of t and the real parameter λ , namely,

$$P(t, \lambda) = e^{-\int_1^t \frac{\gamma}{m(s)} ds} \left(\lambda - g e^{-\int_1^t \frac{\gamma}{m(s)} ds} \int_1^t e^{\int_1^u \frac{\gamma}{m(s)} ds} m(u) du \right),$$

and obtain a complete solution of the Hamilton–Jacobi problem for H :

$$\phi_\lambda : (t, q, z) \longmapsto (t, q, p \equiv P(t, \lambda), z) .$$

Damped forced harmonic oscillator I

- Consider the product manifold $\mathbb{R} \times \mathrm{T}Q \times \mathbb{R}$ with Hamiltonian function

$$H(t, q, p, s) = \frac{p^2}{2m} + \frac{k}{2}q^2 - qF(t) + \frac{\gamma}{m}s.$$

- Conserved quantity:

$$g(t, q, p) = e^{\frac{\gamma t}{2m}} \left(\frac{\sinh\left(\frac{\kappa t}{2m}\right) (2kmq + \gamma p)}{\kappa} + p \cosh\left(\frac{\kappa t}{2m}\right) \right) - \int_1^t F(s) e^{\frac{\gamma s}{2m}} \left(\cosh\left(\frac{\kappa s}{2m}\right) + \frac{\gamma \sinh\left(\frac{\kappa s}{2m}\right)}{\kappa} \right) ds,$$

where $\kappa = \sqrt{\gamma^2 - 4km}$.

Damped forced harmonic oscillator II

- Thus, we can write p in terms of t, q and a real parameter λ as

$$P(t, q, \lambda) = \frac{e^{-\frac{\gamma t}{2m}}}{\gamma \sinh\left(\frac{\kappa t}{2m}\right) + \kappa \cosh\left(\frac{\kappa t}{2m}\right)} \left[\kappa \lambda - 2kmqe^{\frac{\gamma t}{2m}} \sinh\left(\frac{\kappa t}{2m}\right) \right. \\ \left. \kappa \int_1^t e^{\frac{s\gamma}{2m}} F(s) \left(\cosh\left(\frac{\kappa s}{2m}\right) + \frac{\gamma \sinh\left(\frac{\kappa s}{2m}\right)}{\kappa} \right) ds \right],$$

and obtain a complete solution of the Hamilton–Jacobi problem:

$$\Phi_\lambda: (t, q, \lambda, z) \mapsto (t, q, p \equiv P(t, q, \lambda), z).$$

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Thank you!

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